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## **ANALYSIS OF FLUIDIZED BEDS FOR THE SIMULTANEOUS AEROSOL SEPARATION AND HEAT RECOVERY**

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### **ABSTRACT**

A mathematical model is developed to describe the performance of fluidized beds for the simultaneous heat recovery and aerosol separation. This new concept is analyzed in light of the various transport processes taking place within the bed. A two-phase model is developed for the system in which heat and aerosol particles are transferred from the bubble phase to the emulsion phase. In addition to aerosol separation via diffusion, interception, impaction and electrostatic precipitation, thermophoretic collection is also analyzed. The results indicate that high thermal and separation efficiencies can be obtained.

### **INTRODUCTION**

One of the primary sources of atmospheric pollution is the discharge of particulate-laden gases from chemical process industries. Owing to the recent strong emphasis on regulating industrial emissions of aerosols, there has been an increasing trend towards aerosol separation from industrial gas streams prior to their atmospheric disposal. Among the most efficient devices for aerosol

separation are fluidized-bed filters in which a fluidized solid medium is used to capture the particulates.

Fluidized beds have also been used efficiently to recover energy from hot streams. These devices are called fluidized-bed heat regenerators "FBHR's". The essence of FBHR operation is the temporary storage of heat in a fluidized-solid medium.

During the past few years, several research efforts have lead to a reasonable understanding of the performance of FBHR's (1,2) and fluidized-bed aerosol filters (3). As a result of the various transport phenomena taking place when a hot particulate-laden gas is used to fluidize a solid medium, it is not difficult to envision that the fluidized bed can be designed to act as an aerosol filter and a heat regenerator simultaneously. To date, this proposed inherent characteristic of fluidized systems has not been exploited. This new concept can lead to substantial energy recovery, waste reduction and cost savings. The present work aims at developing a mathematical model that can be used to analyze the simultaneous phenomena of heat regeneration and aerosol separation in fluidized beds. The bed performance is investigated via a two-phase model which incorporates various transport processes occurring within the fluidized systems.

### THEORETICAL ANALYSIS

In developing a mathematical model for the system, the following simplifying assumptions are invoked:

1. The bed can be represented by two phases; a bubble phase and an emulsion phase. Flowrate and voidage of the emulsion phase are taken as those at the minimum fluidization condition. This assumption is based on the two-phase model of Toomey and Johnstone (4) which is valid for many fluidization applications.
2. The deposition of the aerosol particles on the solid collectors does not alter the thermal properties of the collectors. This assumption is valid for

collectors whose surface areas are much larger than those of the aerosol particles and for relatively short cycle periods (few hours).

3. Physical properties of the two phases do not change spatially or dynamically. In addition, the gas is assumed to behave ideally such that its density is inversely proportional to temperatures. Therefore, the dynamics of the gas enthalpy can be neglected.
4. The bed operates adiabatically.
5. The bubbles are assumed to be uniform in size throughout the bed. Although the bubbles grow via coalescence as they rise, the assumption of a single average bubble diameter has been successfully used in modeling many fluidization applications (5).

According to the second assumption, the heat-recovery problem can be modeled independently from the aerosol separation phenomenon. In this context, the two-phase heat-regeneration model of El-Halwagi et al. (1) can be adopted. This model assumes that the bubble phase is in plug flow while the emulsion phase is completely mixed. Hence, heat balance around the emulsion phase and an element of the bubble phase can be written, respectively, as

$$W_e C_g (T_o - T_e) + H_{be} e_b A \int_o^H (T_b - T_e) dZ = M_s C_s \frac{dT_e}{dt} \quad (1)$$

$$W_b C_g dT_b + H_{be} e_b A (T_b - T_e) dZ = 0. \quad (2)$$

Equation (2) can be integrated to

$$T_b - T_e = (T_o - T_e) \exp(-m_1 x). \quad (3)$$

where

$$m_1 = \frac{H_{be} e_b A H}{W_b C_g} \quad (4)$$

and

$$x = Z/H \quad (5)$$

Substituting from Eq. (3) into Eq. (1), integrating and rearranging, we get

$$T_e = T_o - (T_o - T_i)\exp(-m_2 t) \quad (6)$$

where

$$m_2 = \frac{[W_e + W_b(1 - \exp(-m_1))]C_g}{M_s C_s} \quad (7)$$

Equation (6) describes the emulsion-phase dynamics. It can be substituted into Eq. (3) to yield

$$T_b = T_o - (T_o - T_i)[1 - \exp(-m_1 x)]\exp(-m_2 t) \quad (8)$$

which presents the transient behavior of the bubble-phase temperature profile.

The average gas temperature at height  $x$  is defined as the temperature that one would measure by sampling the gas at height  $x$  in proportion to the gas flow-rates through the two phases. Hence, Eqs. (6) and (8) can be applied at the bed outlet and combined to give the outlet gas temperature as

$$T_{gH} = T_o - (T_o - T_i) \frac{M_s C_s m_2}{WC_g} \exp(-m_2 t). \quad (9)$$

Having presented the dynamics of the outlet-gas temperature, one can next derive an expression for the thermal-recovery efficiency. The efficiency of heat recovery from a hot gas during a certain time period,  $t_r$ , may be defined as (6):

$$\begin{aligned} \eta_r &= \frac{\text{heat taken up by cold solids in time } t_r}{\text{maximum possible heat gain in time } t_r} \\ &= \frac{\text{heat lost by hot gas in time } t_r}{\text{maximum possible heat loss in time } t_r} \\ &= Q_r / Q_{r_{\max}} \end{aligned} \quad (10)$$

The heat lost by gas in time  $t_r$  is given by

$$Q_r = \int_0^{t_r} WC_g(T_o - T_{gH})dt. \quad (11)$$

Substituting for  $T_{gH}$  from Eq. (9) into Eq. (11) and integrating, we obtain

$$Q_r = M_s C_s (T_o - T_i) [1 - \exp(-m_2 t_r)]. \quad (12)$$

The maximum possible heat transferred in time  $t_r$  may be expressed as

$$Q_{r_{\max}} = WC_g (T_o - T_i) t_r. \quad (13)$$

When Eqs. (12) and (13) are substituted into Eq. (10), we get

$$\eta_r = \frac{M_s C_s \left[ 1 - \exp(-m_2 t_r) \right]}{WC_g \left[ \frac{t_r}{t_r} \right]} \quad (14)$$

which provides a description of the dynamic variation of thermal-recovery efficiency with time.

Having addressed the heat-recovery problem, we are now in a position to tackle the aerosol-separation phenomenon. Owing to the axial and dynamic variations of temperature, a numerical algorithm is needed. Hence, the bed is divided into a number of compartments (stages) in series. Each stage consists of an emulsion phase and a bubble phase; each of which is assumed to be completely mixed. The height of each stage is taken as the average bubble diameter of the bed. Therefore, unlike the heat-transfer model, the mixing patterns for aerosol separation will be described via a stagewise model. This model, which is equivalent to an axial dispersion representation, is appropriate for describing aerosol separation (3). In addition to bubble-emulsion interphase transfer of the aerosol particles via convection/diffusion (7), they can be also transported via thermophoresis induced by the temperature gradient between the two phases. Furthermore, aerosol collection within the emulsion phase can be accomplished

through Brownian diffusion, direct interception, inertial impaction and electrostatic forces. As has been indicated in modeling the heat-transfer operation, the bubble-emulsion temperature gradient changes dynamically. Hence, the rate of interphase transfer of aerosols due to thermophoresis also changes dynamically. Therefore, material balance equations ought to be applied for each time discretely. At a given time  $t$ , material balance equations for the aerosol particles around the  $n^{\text{th}}$  stage of the bubble phase and the emulsion phase are given, respectively, by

$$U_b C_{b,n-1} - U_b C_{b,n} - K_{be} \epsilon_b (C_{b,n} - C_{e,n}) \Delta Z - K_t C_{b,n} = 0 \quad (15)$$

and

$$U_{mf} C_{e,n-1} - U_{mf} C_{e,n} + K_{be} \epsilon_b (C_{b,n} - C_{e,n}) \Delta Z + K_t C_{b,n} - U_{mf} \eta_e \left[ \frac{1.5(1-\epsilon_{mf})}{d_c} \right] (1-\epsilon_b) \Delta Z C_{e,n} = 0. \quad (16)$$

The following boundary conditions can be used:

$$C_{b,o} = C_o \quad (17)$$

and

$$C_{e,o} = C_o. \quad (18)$$

Equations (15) and (16) can be solved simultaneously for each stage to yield  $C_{b,n}$  and  $C_{e,n}$ . The average aerosol concentration leaving the bed is given by

$$\overline{C_N} = (U_b C_{b,N} + U_{mf} C_{e,N})/U \quad (19)$$

and the overall collection efficiency is expressed as

$$E = 1 - \frac{\overline{C_N}}{C_o}. \quad (20)$$

Having derived a complete model for the system, it is necessary to outline expressions for evaluating the model parameters. These are given in the Appendix. For a given operating time, the solution procedure is carried out as follows. First, Eqs. (6) and (8) are used to evaluate the temperatures of the bubble and the emulsion phases throughout the bed. Equation (14) is then employed to evaluate the thermal efficiency of the bed. Having identified temperature gradients throughout the bed, Eqs. (15) and (16) are solved simultaneously for each stage to calculate the concentration of aerosols in the bubble and the emulsion phases. The computation is carried out till the uppermost stage is reached. Hence, Eq. (20) is used to calculate overall collection efficiency. In the sequel, the model is used to investigate the behavior of the proposed system for simultaneous heat recovery and aerosol separation.

### **RESULTS AND DISCUSSION**

In order to analyze the performance of fluidized beds for the simultaneous aerosol separation and heat recovery, the proposed model is used to estimate the dynamic performance of the system. The following data are employed:  $T_i = 300$  K,  $T_o = 800$  K,  $U = 0.3$  m/s,  $U_{mf} = 0.2$  m/s,  $H_{mf} = 2.0$  m,  $D_R = 1.0$  m,  $D_c = 0.002$  m and  $\epsilon_{mf} = 0.45$ . Figure 1 shows the unsteady-state variation in the collector-particle temperature. As the time increases, the temperature of the solid collectors increases. Consequently, both the thermal and the collection efficiencies decrease. Figure 2 shows the dynamic reduction in the bed thermal efficiency. Owing to the continuous reduction in thermal efficiency, the bed ought to be regenerated after a reasonable period of time by temporarily passing a cold fluid instead of the hot gas. Figure 3 illustrates the dynamic variation in the collection efficiency for various aerosol-particle diameters. The nonmonotonic behavior of the aerosol-separation efficiency can be attributed to the shift from a diffusion-limited range to an impaction-limited regime. In this context, fluidized beds offer virtually complete aerosol separation for particle



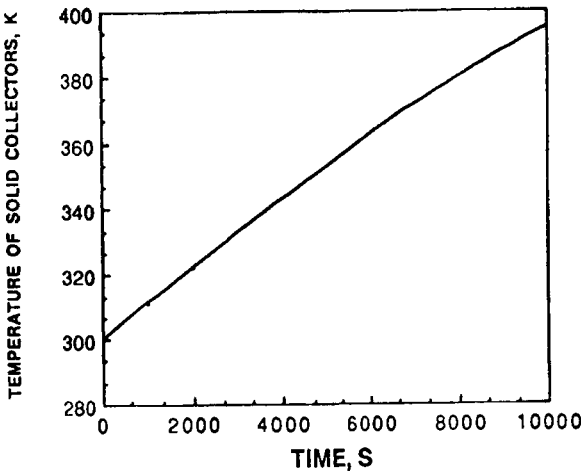


FIGURE 1. Dynamics of solid-collector temperature.

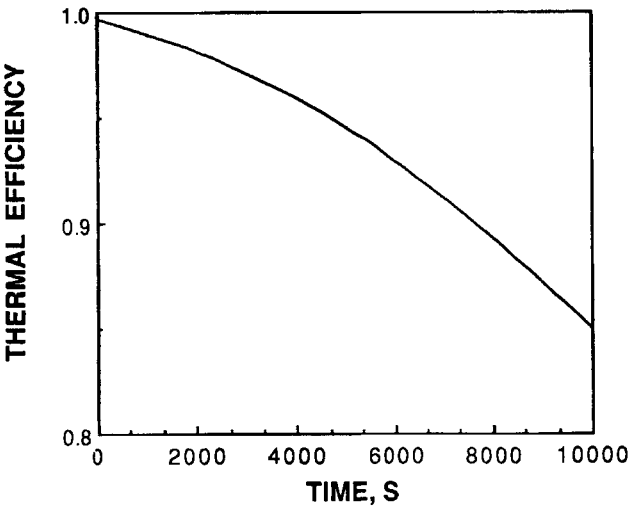
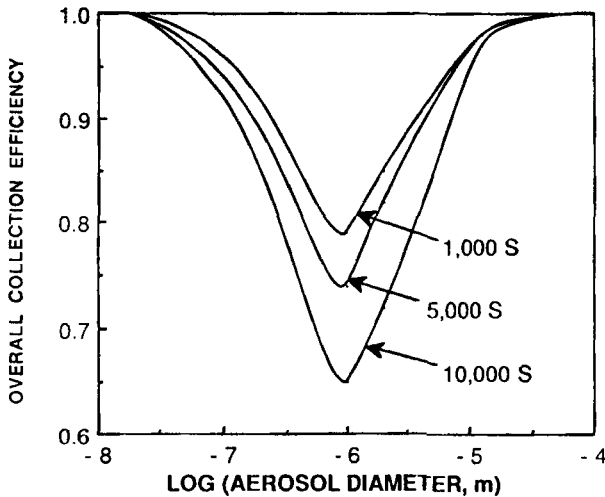


FIGURE 2. Transient thermal efficiency of the bed.



*Fig. 3 Dynamics of Aerosol Collection Efficiency.*

diameters smaller than  $0.01 \mu\text{m}$  and larger than  $100.00 \mu\text{m}$ . In the intermediate range, the collection efficiency becomes less. As time proceeds, the bubble-emulsion thermophoretic temperature gradient decreases. This leads to dynamic reduction in aerosol-separation efficiency.

### CONCLUSION

A mathematical model has been developed to simulate the performance of fluidized beds for the simultaneous aerosol separation and heat recovery. A two-phase model was proposed to describe the heat- and mass-transfer phenomena throughout the bed. The model accounts for separation enhancement due to thermophoresis. The model results indicate that indeed high thermal and aerosol-separation efficiencies can be accomplished in a fluidized bed.

## APPENDIX

### EVALUATION OF MODEL PARAMETERS

In the following, the expressions employed in estimating the model parameters are presented. All of these parameters are related to readily-obtainable quantities.

#### A.1. Bubble Diameter

The correlation of Mori and Wen (8) is used to evaluate the bubble size. It is given by

$$D_b = D_{bm} - (D_{bm} - D_{bo})\exp(-0.3Z/D_R) \quad (\text{A.1})$$

where

$$D_{bm} = 1.6377[A(U - U_{mf})]^{0.4} \quad (\text{A.2})$$

where  $U$  is the superficial velocity of the fluidizing gas, and

$$D_{bo} = 0.376(U - U_{mf})^2 \quad (\text{A.3})$$

(for porous distributor plates)

or

$$D_{bo} = 0.8716 \left[ \frac{A(U - U_{mf})}{N_D} \right]^{0.4} \quad (\text{A.4})$$

(for perforated distributor plates with number of orifice openings =  $N_D$ ).

The average bubble diameter of the bed is assumed to be the value of  $D_b$  at  $H_{mf}/2$ .

#### A.2. Fractional Volume of Bed Occupied by the Bubble Phase, $\epsilon_b$

$$\epsilon_b = (U - U_{mf})/U_b \quad (\text{A.5})$$

where (7)

$$U_b = U - U_{mf} + U_{br} \tag{A.6}$$

and the absolute velocity of an isolated single rising bubble as suggested by Davidson and Harrison (7) is given by

$$U_{br} = 0.711(gD_b)^{0.5} \tag{A.7}$$

**A.3. The Expanded-Bed Height**

The bed expansion above the height at incipient fluidization is assumed to be entirely due to the presence of bubbles, so

$$H = H_{mf} / (1 - \epsilon_b) \tag{A.8}$$

**A.4. Distribution of Gas Flow Among the Phases**

According to the two-phase theory of fluidization (4):

$$W_e = U_{mf} \rho_g A \tag{A.9}$$

and

$$W_b = (U - U_{mf}) \rho_g A \tag{A.10}$$

or

$$U_b = U - U_{mf} \tag{A.11}$$

**A.5. Interphase Exchange Coefficients**

The interphase heat transfer coefficient between the bubble phase and the emulsion phase may be calculated from the following semi-empirical expression (5):

$$H_{be} = 4.5 \left( \frac{U_{mf} \rho_g C_g}{D_b} \right) + 5.85 \frac{(k_g \rho_g C)^{0.5} g^{0.25}}{D_b^{1.25}} \tag{A.12}$$

In a similar fashion, the bubble-emulsion interphase mass transfer coefficient is given by (7):

$$K_{be} = 4.5 U_{mf} / D_b + 5.85 \left( \frac{D_e^{1/2} g^{1/4}}{D_b^{5/4}} \right) \quad (\text{A.13})$$

The thermophoretic velocity is given by (9):

$$K_t = \frac{0.4 k_g}{(k_g + k_p)P} \nabla T \quad (\text{A.14})$$

The temperature gradient can be assumed to take place linearly across the bubble clouds. Thus within the  $n^{\text{th}}$  stage

$$\nabla T \propto \frac{T_{b,n} - T_e}{\delta_n} \quad (\text{A.15})$$

where  $\delta_n$  is the  $n^{\text{th}}$  stage cloud thickness which can be estimated by (10):

$$\delta_n = D_{b,n} \left[ \left( \frac{U_{br} \epsilon_{mf} / U_{mf}}{U_{br} \epsilon_{mf} / U_{mf} - 1} \right)^{1/3} - 1 \right] \quad (\text{A.16})$$

Clearly Eq. (A.15) results in a conservative estimate for the thermophoretic velocity since the actual boundary layer around the bubble occupies just a fraction of the cloud.

#### A.6. Single-Sphere Collection Efficiency

Aerosol filtration is assumed to take place through four collection mechanisms; Brownian diffusion, direct interception, inertial impaction, and induced electrostatic forces. The overall collection efficiencies for a single collector ( $\eta_e$ ) is assumed to be the sum of the four mechanisms in the corresponding phase, i.e.,

$$\eta_e = \eta_D + \eta_{Int} + \eta_{Imp} + \eta_{Elect.} \quad (\text{A.17})$$

The following expressions are used throughout this work:

Brownian diffusion (11)

$$\eta_D = 3.5 \left( \frac{1 - \alpha^{5/3}}{\beta_H} \right)^{1/3} Pe^{-2/3} \quad (\text{A.18})$$

where, Happel's (12) hydrodynamic factor is given by

$$\beta_H = 1 - 1.5 \alpha^{0.33} + 1.5 \alpha^{1.67} - \alpha^2. \quad (\text{A.19})$$

Direct interception (13)

$$\eta_{Int} = \frac{1}{\beta_K} \{ (1 + I)^2 - 1.5(1 + I) + \alpha[0.2/(1 + I) + 0.5(1 + I)^2 - 0.3(1 + I)] \} \quad (\text{A.20})$$

where Kuwabara's (14) hydrodynamic factor is expressed as

$$\beta_K = 1 - 9/5 \alpha^{.33} - .2 \alpha^2 + \alpha. \quad (\text{A.21})$$

Inertial impaction (15):

$$\eta_{Imp} = Stk_{eff}^{3.55} / (1.67 + Stk_{eff}^{3.55}) \quad (\text{A.22})$$

where

$$Stk_{eff} = Stk \left[ \frac{1 - \alpha^{5/3}}{1 - 1.5\alpha^{1/3} + 1.5\alpha^{5/3} - \alpha^2} + \frac{1.14Re^{1/2}}{e^{1.5}} \right] \quad (\text{A.23})$$

Electrostatic forces (16):

$$n_{Elect.} = \left( \frac{15\pi}{8} K_{IC} \right)^{0.4}$$

where

$$\text{Charged-collector image force } K_{IC} = \frac{(e_D - 1)2CD_p^2 Q_{ac}^2}{(e_D + 2)3uD_c U_{Po} e_o}.$$

### NOMENCLATURE

- |                  |   |
|------------------|---|
| A                | cross-sectional area of bed, m <sup>2</sup>   |
| C                | Cunningham correction factor  |
| C <sub>b,n</sub> | gas concentration of aerosol leaving stage <i>n</i> of bubble phase, #/m <sup>3</sup> |

$C_{c,n}$	gas concentration of aerosol leaving stage $n$ of emulsion phase, $\#/m^3$
$C_g$	gas specific heat, J/kg K
$C_o$	inlet aerosol concentration, $\#/m^3$
$\bar{C}_N$	average aerosol concentration leaving the filter, $\#/m^3$
$C_s$	solid-particle specific heat, J/kg K
$D_b$	average equivalent bubble diameter, m
$D_{bm}$	maximum bubble diameter, m
$D_{b,n}$	bubble size within the $n^{\text{th}}$ stage, m
$D_{bo}$	initial bubble diameter, m
$D_c$	diameter of collector particles, m
$D_e$	molecular diffusion coefficient of aerosol particulates, $m^2/s$
$D_p$	diameter of aerosol particulates, m
$D_R$	bed diameter, m
$E$	overall aerosol-collection efficiency of the bed
$g$	acceleration of gravity, $m^2/s$
$H$	bed height, m
$H_{be}$	interphase heat exchange coefficient between bubble phase and emulsion phase (rate of heat exchanged between the bubble phase and the emulsion phase per unit time, per unit volume of bubble phase, per unit temperature difference between the bubble phase and the emulsion phase), $W/m^3 K$
$H_{mf}$	bed height at minimum fluidization conditions, m
$I$	interception number ( $D_p/D_c$ )
$K_{be}$	volumetric rate of gas exchange between bubble and emulsion phases, $s^{-1}$

$k_g$	thermal conductivity of the gas, W/m K
$K_{ic}$	dimensionless characteristic mobility for charged-collector image force
$k_p$	thermal conductivity of the aerosol particles, $\frac{W}{mK}$
$K_t$	thermophoretic velocity, m/s
$m_1$	dimensionless parameter given by Eq. (4)
$m_2$	defined by Eq. (7), $s^{-1}$
$M_s$	total mass of solid particles inside the bed, kg
$N$	total number of stages in bed
$N_D$	number of orifice openings on the distributor
$Q_r$	heat lost by gas or heat uptaken by solids, J
$Q_{r_{max}}$	maximum possible heat lost by hot gas during time $t_r$ or maximum possible heat gained by cold solids during time $t_r$ , J
$Stk$	Stokes number
$Stk_{eff}$	effective Stokes number
$t$	time, s
$t_r$	time during which heat is recovered from the gas by the bed, s
$T_b$	temperature of bubble phase, K
$T_{bH}$	outlet temperature of bubble phase, K
$T_{b,n}$	bubble-phase temperature within the $n^{th}$ stage, K
$T_e$	temperature of emulsion phase, K
$T_g$	average gas temperature at height $Z$ , K
$T_{gH}$	outlet gas temperature, K



$T_i$	initial bed temperature, K
$T_o$	inlet gas temperature, K
$U_b$	absolute rise velocity of a swarm of bubbles through a fluidized bed, m/s
$u_{br}$	absolute rise velocity of single bubble through a fluidized bed, m/s
$U$	superficial velocity of gas, m/s
$U_b$	superficial velocity of the bubble phase, m/s
$U_{mf}$	superficial velocity of gas at minimum fluidization conditions, m/s
$W$	mass flow-rate of the gas, kg/s
$W_b$	mass flow-rate of bubble phase, kg/s
$W_e$	mass flow-rate of emulsion phase, kg/s
$x$	dimensionless fractional height above the distributor
$Z$	height above the distributor, m

### GREEK SYMBOLS

$\alpha$	solidity
$\beta_H$	Happel's hydrodynamic factor, Eq. (A.19)
$\beta_K$	Kuwabara's hydrodynamic factor, Eq. (A.21)
$\delta_a$	cloud thickness within the nth stage
$\Delta Z$	stage height, m
$\epsilon$	fractional voidage
$\epsilon_b$	volume fraction of bubbles in the bed
$\epsilon_D$	dielectric constant of aerosol particulates

$\epsilon_{mf}$	mean voidage under minimum-fluidization conditions
$\epsilon_0$	permittivity of free space, $8.85 \times 10^{-12} \text{ C}^2/\text{N, m}^2$
$\eta_D$	single spherical collection efficiency for Brownian diffusion
$\eta_e$	collection efficiency of a single sphere in the emulsion phase
$\eta_{Elect}$	single spherical collection efficiency for induced electrostatic collection
$\eta_{imp}$	single spherical collection efficiency for impaction
$\eta_{int}$	single spherical collection efficiency for interception
$\eta_r$	thermal efficiency of the regenerator
$\rho_g$	gas density, $\text{kg/m}^3$
$\rho_s$	solid density, $\text{kg/m}^3$

### Subscripts

<b>b</b>	bubble phase
<b>e</b>	emulsion phase
<b>n</b>	stage number
<b>N</b>	final stage

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